Flaglets for studying the largescale structure of the Universe

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Roadmap

- Galaxy surveys & data on the ball
- Fourier-Laguerre transform on the ball
- Flaglet transform on the ball
- Spin directional wavelets on the sphere & ball

Cosmology with galaxy surveys



- Dark matter =invisible
 - Nature and properties of dark matter, dark energy?
 - GR or modified gravity?
 - Origin of structure? Inflation?
 - Signatures imprinted in the large-scale structure

The large-scale structure of the Universe



Angle on the sky +
 redshift = **3D position**

3D cosmic web:
 filaments, walls, voids
 due to hierarchical
 structure formation

Exploiting LSS data *is* complicated

- Photometric surveys (DES, Euclid, LSST): new challenges
- Photo-z errors, spatially varying systematics / depth
- Complicated geometry / selection functions

dust



Systematics is the new frontier

Appropriate methods are essential



Wish list for novel 3D transforms



3D spherical measure

 $\mathrm{d}^3 \vec{r} = r^2 \sin \theta \mathrm{d} \theta \mathrm{d} \phi \mathrm{d} r$

- Separable (data on $\mathbb{S}^2 \times \mathbb{R}^+$ rather than \mathbb{R}^3)
- Meaningful translation, rotation operators
- Theoretically exact / sampling theorem
- Relate to Fourier-Bessel



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Exact spherical harmonic transform

- Spherical harmonics: $f(\theta, \phi) = \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} f_{\ell m} \left[Y_{\ell m}(\theta, \phi) \right]$
- MW sampling theorem : band-limited at L \iff information captured in $N_{\rm pix}\sim 2L^2$ samples



⇒ integrals discretised without any approximation ⇒ theoretically exact transform

McEwen & Wiaux (2011)

Exact spherical Laguerre transform

• Basis functions on
$$\mathbb{R}^+$$

 $K_p(r) \equiv \sqrt{\frac{p!}{(p+2)!}} \frac{e^{-r/2\tau}}{\sqrt{\tau^3}} L_p^{(2)}\left(\frac{r}{\tau}\right)$
• Exact transform:
 $f(r) = \sum_{p=0}^{P-1} f_p K_p(r)$
 $f(r) = \sum_{p=0}^{P-1} f_p K_p(r)$
 $f_p = \sum_{i=0}^{P-1} w_i f(r_i) K_p(r_i)$
• P samples on [0, R]
Leistedt & McEwen (2012)

The Fourier-Laguerre transform

• Basis on $\mathbb{B}^3 = \mathbb{S}^2 \times \mathbb{R}^+$ with measure $d^3 \vec{r} = r^2 \sin \theta d\theta d\phi dr$

$$Z_{\ell m p}(\vec{r}) = K_p(r) Y_{\ell m}(\theta, \phi) \qquad \vec{r} = (r, \theta, \phi)$$

For band limited signals,

 $f_{\ell m p} = 0, \, \forall \ell \ge L, \, \forall p \ge P$

f reconstructed on $\left\{ \begin{array}{l} f_{\ell m p} \end{array} \right\} \sim 2 P L^2$ samples



Leistedt & McEwen (2012)

Connection to Fourier-Bessel analysis

- Basis functions: $Z'_{\ell m}(k; \vec{r}) = Y_{\ell m}(\theta, \phi) j_{\ell}(kr)$
- No sampling theorem for $\int_{\mathbb{R}^+} f(r) j_\ell(kr) r^2 \mathrm{d}r$
- BUT can be calculated (exactly!) from Fourier-Laguerre:

$$\begin{split} \tilde{f}_{\ell m}(k) &= \sqrt{\frac{2}{\pi}} \sum_{p} f_{\ell m p} j_{\ell p}(k) \\ \text{finite sum if} & \text{has analytical} \\ \text{band-limited} & \text{formula} \end{split}$$

Exact evaluation of Fourier-Bessel transform!

Leistedt & McEwen (2012)



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Flaglet transform

Projection:

$$\frac{W^{\Phi}(\vec{r}) \equiv (f \star \Phi)(\vec{r}) = \langle f | \mathcal{T}_{r} \mathcal{R}_{(\theta,\phi)} \Phi \rangle}{W^{\Psi^{jj'}}(\vec{r}) \equiv (f \star \Psi^{jj'})(\vec{r}) = \langle f | \mathcal{T}_{r} \mathcal{R}_{(\theta,\phi)} \Psi^{jj'} \rangle}$$
Reconstruction:

$$\frac{f(r,\theta,\phi) = \int_{\mathbb{B}^{3}} W^{\Phi}(\vec{s})(\mathcal{T}_{r} \mathcal{R}_{(\theta,\phi)} \Phi)(\vec{s}) \mathrm{d}^{3} \vec{s}}{+\sum_{jj'} \int_{\mathbb{B}^{3}} W^{\Psi^{jj'}}(\vec{s})(\mathcal{T}_{r} \mathcal{R}_{(\theta,\phi)} \Psi^{jj'})(\vec{s}) \mathrm{d}^{3} \vec{s}}$$

- Uses translation, rotation, convolution operators on \mathbb{B}^3
- Done in harmonic space thanks to sampling theorem

Leistedt & McEwen (2012)

Convolution on the ball

- 3D operator: $\mathcal{T}_{\vec{r}} \equiv \mathcal{T}_r \mathcal{R}_{(\theta,\phi)}$ radial translation + SO(2) rotation
- Convolution on the ball:

$$(f \star \Psi^{jj'})(\vec{r}) = \langle f | \mathcal{T}_{\vec{r}} \rangle_{\mathbb{B}^3}$$
$$= \int_{\mathbb{B}^3} f(\vec{s}) (\mathcal{T}_{\vec{r}} \Psi^{jj'})(\vec{s}) d\vec{s}$$

Axisymmetric wavelets:

$$(f \star \Psi^{jj'})_{\ell m p} = \sqrt{\frac{4\pi}{2\ell+1}} f_{\ell m p} \Psi^*_{\ell 0 p}$$



Translated flaglets

Scale-discretised wavelets on the sphere



Tiling of the Fourier-Laguerre space

Defining

Tilling of the radial harmonic line







Flaglet transform of Horizon simulation



Underdensities in the Horizon simulation

Ongoing work





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Steerable scale-discretised wavelets

• Change
$$\Psi^j_\ell = \kappa^j(\ell) \implies$$

$$\Psi_{\ell m}^j = \kappa^j(\ell) s_{\ell m}$$

Azimuthal band-limit N

$$s_{\ell m} = 0 \quad \forall \ell, m \text{ with } |m| \ge N$$

Resulting wavelets:



(Wiaux et al 2009, McEwen et al 2013)

Application to Earth tomography data



vvavelet scale j=1, n=3





Use SO3 sampling theorem & exact Wigner transform

Spin directional wavelets on the sphere

- Same wavelets tiling but spin harmonic transforms ${}_sY_{\ell m}$



E-B separation through wavelet transform:

<u>Step 1</u> : forward **spin** wavelet transform

<u>Step 2</u> : inverse **scalar** wavelet transform

$$(Q+iU)(\theta,\phi) \longrightarrow \{ W^{2}\Psi_{j}(\theta,\phi,\rho) \}$$

Real $\{W^{2}\Psi_{j}(\theta,\phi,\rho)\} \longrightarrow -\tilde{E}(\theta,\phi)$
Imag $\{W^{2}\Psi_{j}(\theta,\phi,\rho)\} \longrightarrow -\tilde{B}(\theta,\phi)$

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Application: wavelet denoising of spin data



Spin signal (Q) - Denoised map

Spin signal (U) - Denoised map



Spin directional **Flaglets**

Directional flaglets

Ongoing work



Spin 2 directional flaglets



Application to 3D weak lensing analysis (e.g., E-B sep)



Summary & conclusions



- Exact Fourier-Laguerre and Flaglet transforms
- Exact Wigner and spin directional wavelet transforms
- <u>Ongoing</u>: E-B separation, spin directional flaglets
- Future: complex galaxy survey data, cosmic voids, CMB E-B separation, 3D weak lensing, etc

<u>www.flaglets.org</u> <u>www.s2let.org</u>

1205.0792, 1308.5480, 1308.5406, + papers in prep

Extra Slides

Flaglet denoising of geophysical model

S40RTS: Ritsema's seismological Earth model of shear wavespeed perturbations in the mantle



Flaglets $\Psi^{jj'}$



The spherical Laguerre sampling theorem

Spherical Laguerre basis:

$$K_p(r) \equiv \sqrt{\frac{p!}{(p+2)!}} \frac{e^{-r/2\tau}}{\sqrt{\tau^3}} L_p^{(2)}\left(\frac{r}{\tau}\right)$$



 f band-limited at P : projected/ reconstructed on P samples



- Rescaling on any intervals [0,R]
- Sampling denser near origin due to measure $r^2 \mathrm{d}r$

Scale-discretised generating functions (1)

Smooth generating functions

$$s_{\lambda}(t) \equiv s\left(\frac{2\lambda}{\lambda-1}(t-1/\lambda)-1\right)$$
$$s(t) \equiv \begin{cases} e^{-\frac{1}{1-t^2}}, & t \in [-1,1]\\ 0, & t \notin [-1,1] \end{cases}$$
$$k_{\lambda}(t) \equiv \frac{\int_{t}^{1} \frac{dt'}{t'} s_{\lambda}^{2}(t')}{\int_{1/\lambda}^{1} \frac{dt'}{t'} s_{\lambda}^{2}(t')},$$

Azisymmetric wavelets filters

$$\Psi_{\ell m}^{j} \equiv \sqrt{\frac{2\ell+1}{4\pi}} \kappa_{\lambda} \left(\frac{\ell}{\lambda^{j}}\right) \delta_{m0}.$$
$$\Phi_{\ell m} \equiv \sqrt{\frac{2\ell+1}{4\pi}} \eta_{\lambda} \left(\frac{\ell}{\lambda^{J_{0}}}\right) \delta_{m0}.$$

Tilling of Fourier-Laguerre space

$$\kappa_{\lambda}(t) \equiv \sqrt{k_{\lambda}(t/\lambda) - k_{\lambda}(t)} \quad \text{and} \quad \eta_{\lambda}(t) \equiv \sqrt{k_{\lambda}(t)}$$
$$\kappa_{\nu}(t) \equiv \sqrt{k_{\nu}(t/\nu) - k_{\nu}(t)} \quad \text{and} \quad \eta_{\nu}(t) \equiv \sqrt{k_{\nu}(t)}$$
$$\eta_{\lambda\nu}(t,t') \equiv \sqrt{k_{\lambda}(t/\lambda)k_{\nu}(t') + k_{\lambda}(t)k_{\nu}(t'/\nu) - k_{\lambda}(t)k_{\nu}(t')}.$$

$$\underbrace{\operatorname{let}}_{\ell m p} = \sqrt{\frac{2\ell+1}{4\pi}} \kappa_{\lambda} \left(\frac{\ell}{\lambda^{j}}\right) \kappa_{\nu} \left(\frac{p}{\nu^{j'}}\right) \delta_{m0}.$$

$$\underbrace{\operatorname{rs:}}_{\Phi_{\ell m p}} = \begin{cases} \sqrt{\frac{2\ell+1}{4\pi}} \eta_{\nu} \left(\frac{p}{\nu^{J_{0}'}}\right) \delta_{m0}, & \text{if } \ell > \lambda^{J_{0}}, \ p \le \nu^{J_{0}'} \\ \sqrt{\frac{2\ell+1}{4\pi}} \eta_{\lambda} \left(\frac{\ell}{\lambda^{J_{0}}}\right) \delta_{m0}, & \text{if } \ell \le \lambda^{J_{0}}, \ p > \nu^{J_{0}'} \\ \sqrt{\frac{2\ell+1}{4\pi}} \eta_{\lambda\nu} \left(\frac{\ell}{\lambda^{J_{0}}}, \frac{p}{\nu^{J_{0}'}}\right) \delta_{m0}, & \text{if } \ell < \lambda^{J_{0}}, \ p < \nu^{J_{0}'} \\ 0, & \text{elsewhere.} \end{cases}$$

Fourier-Bessel transform

- Fourier-Bessel : $Z'_{\ell m}(k; \vec{r}) = Y_{\ell m}(\theta, \phi) j_{\ell}(kr)$
- Eigenfunctions of the Laplacian in 3D spherical coord.

$$f(\vec{r}) = \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} \sqrt{\frac{2}{\pi}} \int_{\mathbb{R}^{+}} \mathrm{d}k k^{2} \tilde{f}_{\ell m}(k) Y_{\ell m}(\theta,\phi) j_{\ell}(kr)$$

$$\tilde{f}_{\ell m}(k) = \sqrt{\frac{2}{\pi}} \int_{\mathbb{S}^{2}} \mathrm{d}\Omega(\theta,\phi) \int_{\mathbb{R}^{+}} \mathrm{d}rr^{2} f(r,\theta,\phi) Y_{\ell m}^{*}(\theta,\phi) j_{\ell}(kr)$$

• But no sampling theorem for $\int_{\mathbb{R}^+} f(r) j_\ell(kr) r^2 dr$

Connection to Fourier-Bessel analysis (2)

$$\begin{split} \tilde{f}_{\ell m}(k) &= \sqrt{\frac{2}{\pi}} \sum_{p} f_{\ell m p} j_{\ell p}(k) \quad \text{with} \quad j_{\ell p}(k) \equiv \langle K_{p} | j_{\ell} \rangle \\ & \underline{\text{finite sum if band-limited}} & \underline{\text{has analytical formula}} \\ j_{\ell p}(k) &= \sqrt{\frac{p!}{(p+2)!}} \sum_{j=0}^{p} c_{j}^{p} \mu_{j+2}^{\ell}(k), \\ & \underline{\text{Details of}} \\ & \text{analytic formula:} \quad c_{j}^{p} \equiv \frac{(-1)^{j}}{j!} \binom{p+2}{p-j} = -\frac{p-j+1}{j(j+2)} c_{j-1}^{p}. \\ & \mu_{j}^{\ell}(k) \equiv \frac{1}{\tau^{j-\frac{1}{2}}} \int_{\mathbb{R}^{+}} \mathrm{d}rr^{j} j_{\ell}(kr) e^{-\frac{r}{2\tau}} \\ & \text{with} \\ & \mu_{j}^{\ell}(k) = \sqrt{\pi} \ 2^{j} \ \tilde{k}^{\ell} \ \tau^{\frac{3}{2}} \ \frac{\Gamma(j+\ell+1)}{\Gamma(\ell+\frac{3}{2})} \ {}_{2}F_{1}\left(\frac{j+\ell+1}{2}; \frac{j+\ell}{2}+1; \ell+\frac{3}{2}; -4\tilde{k}^{2}\right) \end{split}$$

More details on the Flaglet transform

Thanks to sampling theorem, flaglet transform easily computed

Projection:

$$\begin{aligned}
W^{\Phi}(\vec{r}) &\equiv (f \star \Phi)(\vec{r}) &= \langle f | \mathcal{T}_{r} \mathcal{R}_{(\theta,\phi)} \Phi \rangle \\
W^{\Psi^{jj'}}(\vec{r}) &\equiv (f \star \Psi^{jj'})(\vec{r}) &= \langle f | \mathcal{T}_{r} \mathcal{R}_{(\theta,\phi)} \Psi^{jj'} \rangle \\
f(r,\theta,\phi) &= \int_{\mathbb{B}^{3}} W^{\Phi}(\vec{s}) (\mathcal{T}_{r} \mathcal{R}_{(\theta,\phi)} \Phi)(\vec{s}) \mathrm{d}^{3} \vec{s} \\
&+ \sum_{jj'} \int_{\mathbb{B}^{3}} W^{\Psi^{jj'}}(\vec{s}) (\mathcal{T}_{r} \mathcal{R}_{(\theta,\phi)} \Psi^{jj'})(\vec{s}) \mathrm{d}^{3} \vec{s}
\end{aligned}$$

Most efficient: filtering in Fourier-Laguerre space

$$\begin{aligned} \text{Projection:} \quad & W_{\ell m p}^{\Psi^{j j'}} = \sqrt{\frac{4\pi}{2\ell + 1}} f_{\ell m p} \Psi_{\ell 0 p}^{j j'} \qquad & W_{\ell m p}^{\Phi} = \sqrt{\frac{4\pi}{2\ell + 1}} f_{\ell m p} \Phi_{\ell 0 p} \\ \text{Reconstruction:} \quad & f_{\ell m p} = \sqrt{\frac{4\pi}{2\ell + 1}} W_{\ell m p}^{\Phi} \Phi_{\ell 0 p} + \sqrt{\frac{4\pi}{2\ell + 1}} \sum_{j j'} W_{\ell m p}^{\Psi^{j j'}} \Psi_{\ell 0 p}^{j j'} \end{aligned}$$

E-B separation through wavelet transform

<u>Step 1</u> : forward **spin** wavelet transform

<u>Step 2</u> : inverse **scalar** wavelet transform

$$\begin{aligned} (Q+iU)(\theta,\phi) &\longrightarrow \{ W^{2}\Psi_{j}(\theta,\phi,\rho) \} \\ \operatorname{Real}\{W^{2}\Psi_{j}(\theta,\phi,\rho)\} &\longrightarrow -\tilde{E}(\theta,\phi) \\ \operatorname{Imag}\{W^{2}\Psi_{j}(\theta,\phi,\rho)\} &\longrightarrow -\tilde{B}(\theta,\phi) \end{aligned}$$

